# A linear relation to calculate one particle expectation value <r-1> for $1 S 2$ configuration for $2<Z<100$ <br> WISSAM AHMED AMEEN 

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## ARTICLEINFO

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ABSTRACT
One particle expectation value $\langle\mathrm{r}-1>$ for 1 S 2 configuration for $\mathrm{Z}=2$ to 10 are evaluated by using Hartree-Fock-Roothaan (HFR) wavefunction. A linear relation are determined to describe these results, and this relation have tested to check its ability to calculate one particle expectation value <r-1> for atoms with higher atomic numbers ( $\mathrm{Z}=2$ to 100 ). Results are in good agreement with calculations given by [S.L. Saito Atomic Data and Nuclear Data Tables95, (2009), 836-870] for Z<20.

## Introduction:

The $1 \mathrm{~S}^{2}$ configuration plays an important role in the quantum mechanical calculations for its simplicity and applicability as a system to calculate some of its atomic properties like the evaluation of the one particle expectation value with power $n=-1\left\langle r_{1}{ }^{-1}\right\rangle$, the subject of this research. It's very important to calculate $\left\langle\mathrm{r}_{1}{ }^{-1}\right\rangle$ because of its direct relation with the evaluation of the total energy for the system under studying. Different methods were used in quantum mechanical calculations but Hartree-Fock (HF) equations were first proposed by Fock in 1930 has taken a central role in studies of atomic and molecular electronic structure. Atomic HF equations give physical properties of various atomic systems in the nonrelativistic independent particle approximation. These data are also important for molecular calculation, for it is impossible to develop basis sets and pseudo-potentials for use in molecular calculations without them [1][2].

[^0]But these methods still needs many sophisticated mathematical and programming techniques to evaluate an atomic property. In this research a simple linear relation was determined to evaluate one particle expectation value with power $n=-1$ for $1 \mathrm{~S}^{2}$ configuration without any need to use the traditional procedures to calculate this property. This relation may help as a fast tool to check published atomic data and may be to guess an approximate values for $\left\langle\mathrm{r}_{1}{ }^{-1}\right\rangle$.

## Theory:

One particle expectation value for $\mathrm{n}=-1$ given by the following equation:

$$
<\mathrm{r}_{1}^{-1}>=\int_{0}^{\infty} D\left(r_{1}\right) r_{1}^{-1} d r_{1},
$$

Where
$D(r)$ : is the radial electronic density distribution for an electronic system and given by:

$$
\begin{aligned}
& D\left(r_{1}\right)=\int_{0}^{\pi} \int_{0}^{2 \pi} r_{1}^{2} \rho\left(r_{1}\right) d \Omega \\
& \rho\left(r_{1}\right)=N \int \psi^{*}\left(x_{1}, x_{2}, . ., x_{N}\right) \\
& \times \psi\left(x_{1}, x_{2}, . ., x_{N}\right) d s_{1} d x_{2} \ldots d s_{N} d x_{N}
\end{aligned}
$$

The standard notation $d \Omega=\sin (\theta) d \theta d \phi$ is employed, $x_{i}$ denotes a combined spatial and spin coordinate and $\psi\left(x_{1}, x_{2}, . ., x_{N}\right)$ is normalized [3].

The HF wave function $\psi_{H F}$ are used to perform calculations in this research.

The HF wave function for two electron system is:

$$
\psi_{H F}(1,2)=\frac{1}{\sqrt{2}}\left[\Phi_{1}(1) \Phi_{2}(2)-\Phi_{1}(2) \Phi_{2}(1)\right]
$$

Where

$$
\Phi_{j}(r, \theta, \phi)=R(r) Y(\theta, \phi)
$$

$R(r)$ is one particle radial wavefunction and $Y(\theta, \phi)$ is the one particle angular wavefunction of the total wavefunction $\psi_{H F}(1,2)$.

Slater-Type Orbital's (STO's) [4] are used within HF wavefunction in equation (4), the radial part of STO's given by the following equation:

$$
\begin{equation*}
\mathrm{R}(\mathrm{r})=\sum_{k=1}^{M} C_{k} N_{k} r^{n_{i}-1} e^{-\left(\xi_{k}\right) r_{k}} \tag{5}
\end{equation*}
$$

$M$ : Number of basis sets
$n_{k}$ : Principle quantum number
$\xi_{k}$ : Orbital exponent
$C_{k}$ : Expansion coefficient.
$N_{k}$ : Normalization constant and given by:

$$
N_{k}=\sqrt{\frac{\left(2 \xi_{k}\right)^{2 n_{k}+1}}{\left(2 n_{k}\right)!}}
$$

The final equation that used to calculate the expectation value by using data give by [5] as follow:

$$
\begin{gather*}
<r_{1}^{-1}> \\
=\sum_{k=j}^{M} C_{k} C_{j} N_{k} N_{j} \frac{\left(2 n_{k}-1\right)!}{\left(2 \zeta_{k}\right)^{2 n_{k}}}  \tag{6}\\
+2 \sum_{j<k}^{M} C_{k} C_{j} N_{k} N_{j} \frac{\left(n_{k}+n_{j}-1\right)!}{\left(\zeta_{k}+\zeta_{j}\right)^{n_{i}+n_{j}}}
\end{gather*}
$$

The one particle expectation value with power $n=-1$ increases linearly throughout the series as shown in figure(1), so that, the difference between any two values in the series given by:

$$
<\mathrm{r}_{1}^{-1}>^{\Delta_{i, i+1}=\left\langle r_{1}^{-1}>_{i+1}-<r_{1}^{-1}>_{i}\right.}
$$



Figure(1): Liner increasing of $\left.<\mathrm{r}_{1}{ }^{-1}\right\rangle$ with Z , all results in atomic units (a.u.).

So that, the expectation value $<r_{1}{ }^{-1}>_{i+1}$ for any atom throughout the 2 electron like system for $1 S^{2}$ configuration produced by an additive process, and this phenomenon are used in this work to calculate $\left.<r_{1}{ }^{-1}\right\rangle$ for $Z=2$ to 100 by using the following equation:

$$
<r_{1}^{-1}>_{i+1}=<r_{1}^{-1}>_{i}+\Delta_{i, i+1}
$$

The value of $\Delta_{i, i+1}$ was taken as the average for 10 atoms and equal to 0.99128 .

## Results and Conclusions:

Tables (1) and figure (2) shows the results of one particle expectation value with $n=-1$ that calculated by using data from Ref.[5] within equation (6), for atoms of $Z=2$ to 10 in addition to the results from equation (7) that represents the new approach in this research for $Z=2-100$, all results was compared with results of Ref.[2]. All these calculations have made by using MathCad 14.

The values of $<r_{1}{ }^{-1}>$ as shown in figure (1) can be divided into three zones, the first for the interval of $2<Z<8$ and the second is for the interval $8<Z<16$ and the third is for the interval $16<Z<100$. The first zone has a very small differences with respect to results given by the traditional method that given by equation (6) and listed within [2], but the best results of the new approach of this research that given in equation (7) appears
clearly in the second zone where the differences are very small with respect to published work. The third zone shows that the differences between the results given by equation (7) and that given by [2] increases linearly with the increasing of $Z$ within this zone and this may be happened because the roots of the linear relation in equation (7) that emerged from HF calculations with its well-known limitations for $Z>20$ because of the neglecting of relativistic effects [7]. so that the main conclusion in this research is that the linear relation in this research can be used to calculate the one particle expectation value for power $\mathrm{n}=-1$ for $\mathrm{Z}=2-20$ to give a results in good agreement with the published data for $1 \mathrm{~S}^{2}$ configuration.

Table(1): Expectation values for $\mathrm{Z}=2-10$ by using HFR, expectation value for $\mathrm{Z}=2-100$ by using the common difference approach $\Delta_{\mathrm{i}, \mathrm{i}+1}$, all results in atomic units (a.u.).

| z | $<r_{1}{ }^{-1}$ |  | $<r_{1}{ }^{-1}$ | $<r_{1}{ }^{-1}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $>_{\text {HFR }}^{a}$ | $\Delta_{i, i+1}$ | $>^{a}$ | $>^{\text {b }}$ | $\delta^{c}$ |
| 2 | 1.68729 |  | 1.687290 | 1.687282 | 0.00001 |
| 3 | 2.68386 | 0.996570 | 2.678574 | 2.685034 | 0.00646 |
| 4 | 3.676414 | 0.992554 | 3.669857 | 3.681877 | 0.01202 |
| 5 | 4.668731 | 0.992317 | 4.661141 | 4.674339 | 0.01320 |
| 6 | 5.660406 | 0.991675 | 5.652425 | 5.664439 | 0.01201 |
| 7 | 6.65196 | 0.991554 | 6.643708 | 6.653236 | 0.00953 |
| 8 | 7.637847 | 0.985887 | 7.634992 | 7.642171 | 0.00718 |
| 9 | 8.629412 | 0.991565 | 8.626275 | 8.630362 | 0.00409 |
| 10 | 9.617559 | 0.988147 | 9.617559 | 9.618054 | 0.00050 |


| 10.60884 | 10.60738 | 0.00146 |
| :--- | :--- | :--- |
| 11.60013 | 11.59795 | 0.00217 |
|  |  |  |
| 12.59141 | 12.58925 | 0.00216 |
|  |  |  |
| 13.58269 | 13.58115 | 0.00154 |
|  |  |  |
| 14.57398 | 14.57355 | 0.00042 |
|  |  |  |
| 15.56526 | 15.56645 | 0.00119 |
| 16.55654 | 16.55969 | 0.00315 |
| 17.54783 | 17.55323 | 0.00540 |

$18.53911 \quad 18.547360 .00825$
$19.53040 \quad 19.54193 \quad 0.01154$
$20.52168 \quad 20.53754 \quad 0.01586$
$21.51296 \quad 21.53364 \quad 0.02068$
$22.50425 \quad 22.53012 \quad 0.02587$
$23.49553 \quad 23.52726 \quad 0.03173$
$24.48681 \quad 24.52397 \quad 0.03716$
$25.4781 \quad 25.52120 \quad 0.04310$
$26.46938 \quad 26.518660 .04928$
$27.46066 \quad 27.51630 \quad 0.05563$
$28.45195 \quad 28.51438 \quad 0.06244$
$29.44323 \quad 29.51205 \quad 0.06882$
30.4345230 .509720 .07521
$31.42580 \quad 31.50738 \quad 0.08158$

| 33 | 32.41708 | 32.50504 | 0.08796 | 55 | 54.22532 | 54.46797 | 0.24265 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 34 | 33.40837 | 33.50273 | 0.09436 | 56 | 55.21661 | 55.46669 | 0.25008 |
| 35 | 34.39965 | 34.50043 | 0.10078 | 57 | 56.20789 | 56.46547 | 0.25758 |
| 36 | 35.39093 | 35.49815 | 0.10722 | 58 | 57.19917 | 57.46444 | 0.26526 |
| 37 | 36.38222 | 36.49597 | 0.11375 | 59 | 58.19046 | 58.46358 | 0.27312 |
| 38 | 37.37350 | 37.49385 | 0.12035 | 60 | 59.18174 | 59.46262 | 0.28088 |
| 39 | 38.36478 | 38.49189 | 0.12711 | 61 | 60.17302 | 60.46169 | 0.28867 |
| 40 | 39.35607 | 39.49003 | 0.13396 | 62 | 61.16431 | 61.46080 | 0.29649 |
| 41 | 40.34735 | 40.48833 | 0.14098 | 63 | 62.15559 | 62.45994 | 0.30434 |
| 42 | 41.33864 | 41.48663 | 0.14800 | 64 | 63.14688 | 63.45898 | 0.31211 |
| 43 | 42.32992 | 42.48493 | 0.15501 | 65 | 64.13816 | 64.45827 | 0.32011 |
| 44 | 43.32120 | 43.48343 | 0.16222 | 66 | 65.12944 | 65.45748 | 0.32804 |
| 45 | 44.31249 | 44.48192 | 0.16943 | 67 | 66.12073 | 66.45671 | 0.33599 |
| 46 | 45.30377 | 45.48052 | 0.17675 | 68 | 67.11201 | 67.45597 | 0.34396 |
| 47 | 46.29505 | 46.47906 | 0.18401 | 69 | 68.10329 | 68.45525 | 0.35195 |
| 48 | 47.28634 | 47.47763 | 0.19129 | 70 | 69.09458 | 69.45455 | 0.35997 |
| 49 | 48.27762 | 48.47620 | 0.19858 | 71 | 70.08586 | 70.45377 | 0.36791 |
| 50 | 49.26890 | 49.47479 | 0.20589 | 72 | 71.07714 | 71.4530 | 0.37586 |
| 51 | 50.26019 | 50.47339 | 0.21320 | 73 | 72.06843 | 72.45225 | 0.38382 |
| 52 | 51.25147 | 51.47201 | 0.22054 | 74 | 73.05971 | 73.45151 | 0.39180 |
| 53 | 52.24276 | 52.47064 | 0.22789 | 75 | 74.05100 | 74.45078 | 0.39979 |
| 54 | 53.23404 | 53.46929 | 0.23525 | 76 | 75.04228 | 75.45007 | 0.40779 |

$\left.\begin{array}{|l|l|}\hline 76.03356 & 76.44936\end{array}\right) 0.4158$ ( 77.44869 0.42385
$83.96383 \quad 84.44390 \quad 0.48007$
$84.95511 \quad 85.44322 \quad 0.48811$
$85.94640 \quad 86.442560 .49616$
$86.93768 \quad 87.44190 \quad 0.50422$
$87.92897 \quad \mathbf{8 8 . 4 4 1 2 7} \quad \mathbf{0 . 5 1 2 3 0}$
$88.92025 \quad 89.44064 \quad 0.520390$
$89.91153 \quad 90.44011 \quad 0.52858$
$90.90282 \quad 91.43955 \quad 0.53674$
$91.89410 \quad 92.43901 \quad 0.54491$
$92.88538 \quad 93.43851 \quad 0.55313$
$\begin{array}{lll}93.87667 & 94.43799 & 0.56132\end{array}$
$94.86795 \quad 95.437440 .56949$
$95.85923 \quad 96.436940 .57770$
$96.85052 \quad 97.43648 \quad 0.58597$

| 99 | 97.84180 | 98.4360 | 0.59420 |
| :---: | :---: | :---: | :---: |
| 100 | 98.83309 | 99.43553 | 0.60244 |

## a: Present work

b: Ref.[2]
c: $\delta=\left|<r^{-1}>^{a}-<r^{-1}>^{b}\right|$


Figure(2): The relative difference $\delta$ between the results of linear equation and HFR one for $2 \leq \mathrm{Z} \leq 100$.
Where $\delta=\left|<\boldsymbol{r}^{-1}>^{a}-<\boldsymbol{r}^{-1}>^{b}\right|$, all results in this figure are in atomic units (a.u.).

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علاقة خطية لحساب القيمة المتوقعة لبعد الالكترون عن النواة >r-1> للغلاف 1S2 للأعداد الذريـة $2<Z<100$

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الخلاصة
تم حساب القيمة المتوقعة لجسم واحد >r-1> للغلاف 1S2 لللزرات ذات اللأعداد النرية Z من 2 الى 10 باستخدام دالة هارتري -فوك روثان ، ومن ثم تم ايجاد علاقة خطية تصف هذه البيانات، وقد تم اختبار هذه العلاقة لمعرفة امكانية استخدامها لحساب ذرات ذات اعداد ذرية اكبر (Z Z من 2 الى 100)



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